Weak interaction of hadrons.

- Weak interactions of hadrons: constituent quarks emit or absorb W bosons

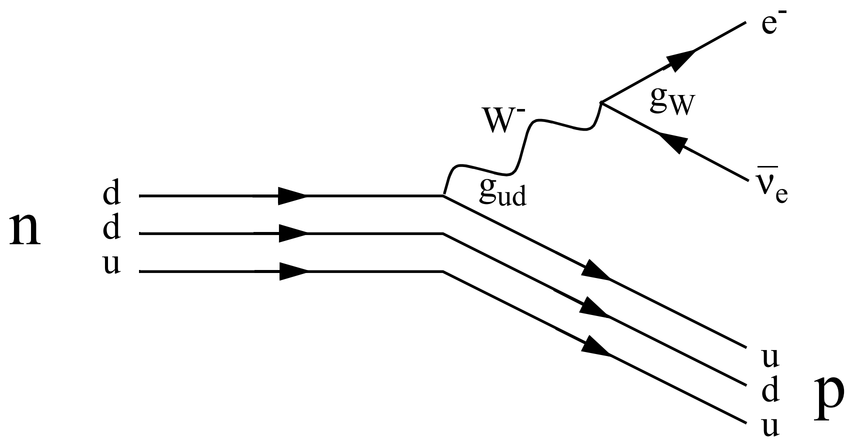


Figure 93: Neutron β-decay.

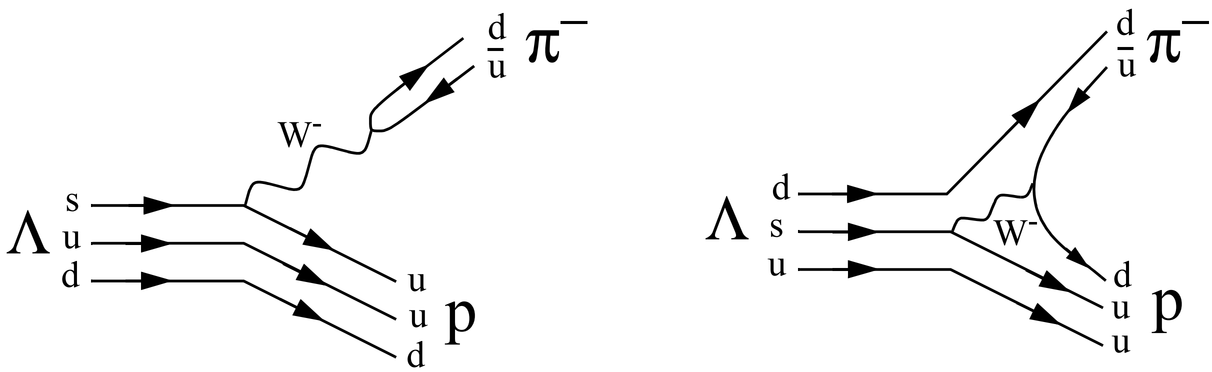


Figure 94: The dominant quark diagrams for Λ decay.

❖ *Lepton-quark symmetry*: corresponding generations of quarks and leptons have identical weak interactions:

$$\begin{pmatrix} \nu_e \\ e^- \end{pmatrix} \leftrightarrow \begin{pmatrix} u \\ d \end{pmatrix} \quad \begin{pmatrix} \nu_\mu \\ \mu^- \end{pmatrix} \leftrightarrow \begin{pmatrix} c \\ s \end{pmatrix} \quad \begin{pmatrix} \nu_\tau \\ \tau^- \end{pmatrix} \leftrightarrow \begin{pmatrix} t \\ b \end{pmatrix}$$

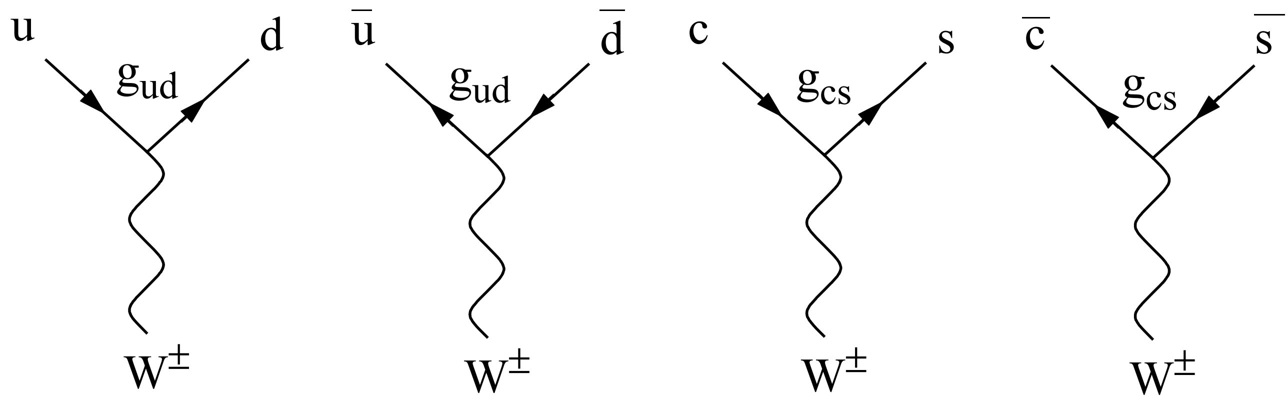


Figure 95: The W-quark vertices obtained from lepton-quark symmetry if quark mixing is ignored.

→ Examples of reactions allowed and not allowed in the lepton-quark symmetry scheme are pion and kaon decay:

$\pi^- \to \mu^- + \bar{\nu}_\mu$	$d\bar{u} \to \mu^- + \bar{\nu}_\mu$	allowed !
$K^- \to \mu^- + \bar{\nu}_\mu$	$s\bar{u} \to \mu^- + \bar{\nu}_\mu$	not allowed !

→ The kaon decay can be made allowed by introducing quark mixing (originally proposed by Cabibbo).

→ According to the quark mixing scheme, d- and s-quarks participate in the weak interactions via the linear combinations:

$$d' = d\cos\theta_C + s\sin\theta_C$$

$$s' = -d\sin\theta_C + s\cos\theta_C$$

where the parameter θ_C is called the *Cabibbo angle*.

→ With quark mixing the quark-lepton symmetry applies to doublets like $\begin{pmatrix} u \\ d' \end{pmatrix}$ *and* $\begin{pmatrix} c \\ s' \end{pmatrix}$

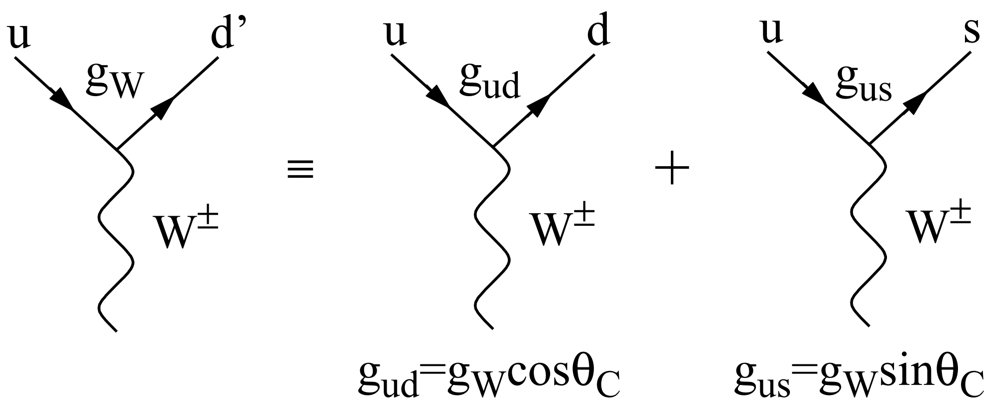


Figure 96: The ud´W vertex can be interpreted as a sum of the udW and usW vertices.

→ With the quark mixing hypothesis some more W-quark vertices are allowed:

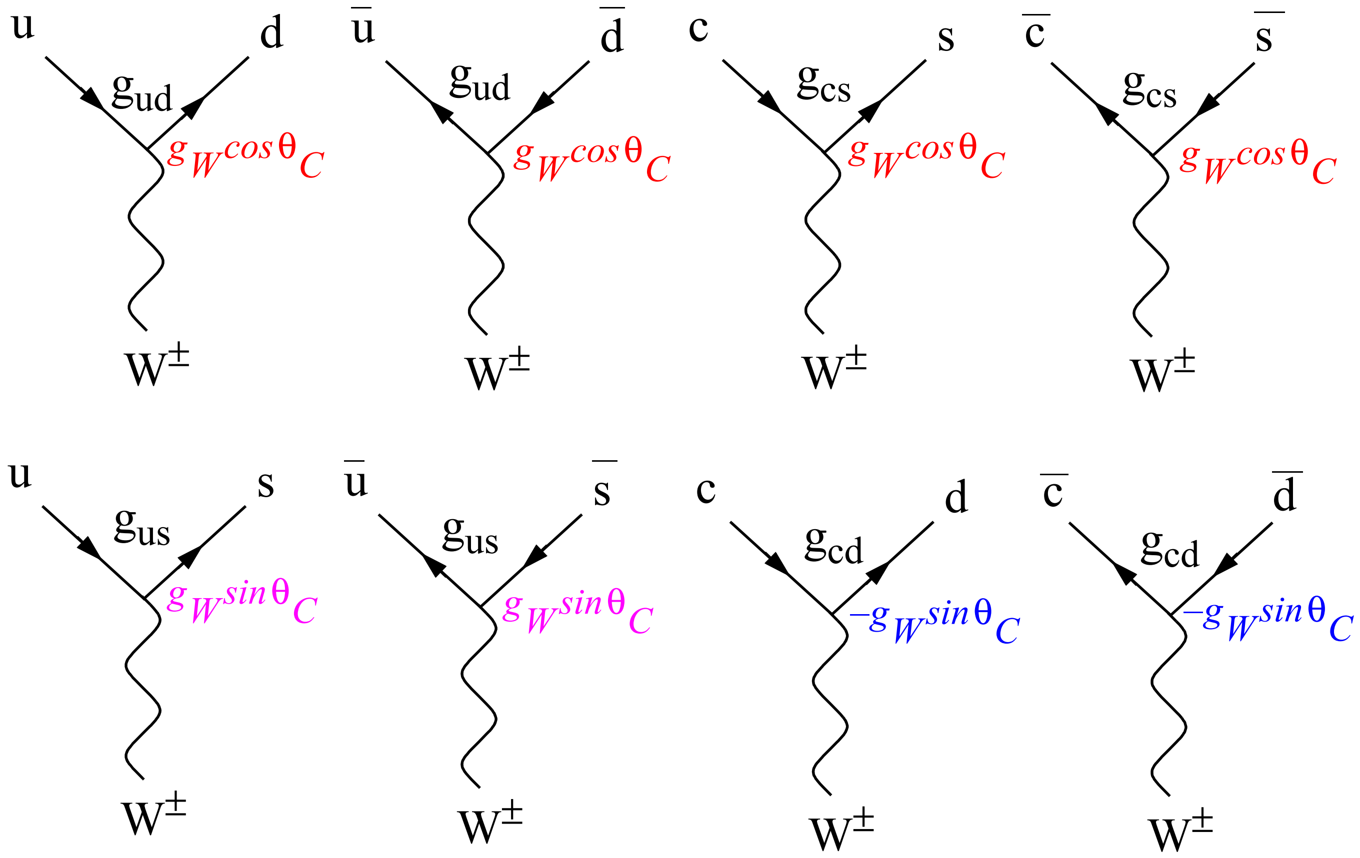


Figure 97: The basic W-quark vertices if quark mixing is taken into account.

The top row of diagrams have couplings given by

$$g_{ud} = g_{cs} = g_W cos\theta_C \qquad (102)$$

while the bottom row diagrams have the couplings

$$g_{us} = -g_{cd} = g_W sin\theta_C \qquad (103)$$

❖ The Cabibbo angle is measured experimentally by comparing decay rates like:

$$\frac{\Gamma(K^- \to \mu^- \bar{\nu}_\mu)}{\Gamma(\pi^- \to \mu^- \bar{\nu}_\mu)} \propto \frac{g_{us}^2}{g_{ud}^2} = \tan^2\theta_C$$

which give the result

$$\theta_C = 12{,}7° \pm 0{,}1° \qquad (104)$$

$$g_W \cos\theta_C = 0{,}98 g_W$$

$$g_W \sin\theta_C = 0{,}22 g_W$$

→ The charmed quark couplings g_{cd} and g_{cs} are measured in neutrino scattering experiments and this give the same result:

$$\theta_C = 12° \pm 1°$$

❖ Experimentally it has been observed that charmed hadrons almost always decays into strange hadrons.

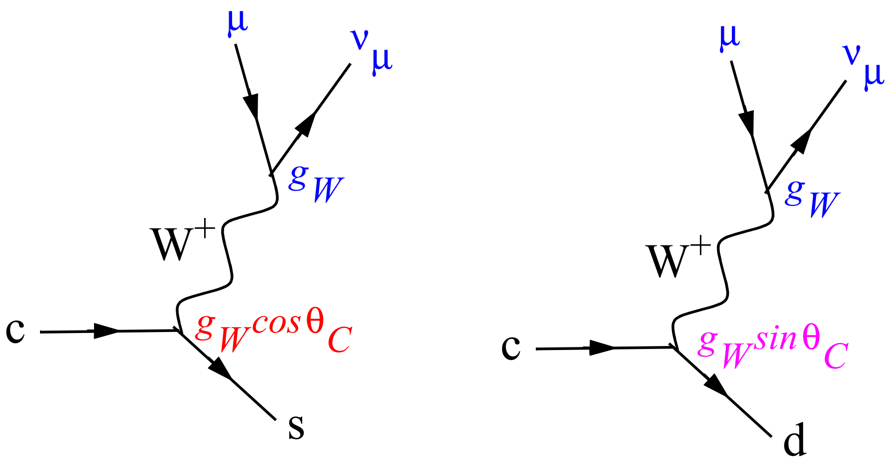


Figure 98: Cabbibo-allowed and Cabbibo-suppressed semi-leptonic decay of a charmed quark.

→ The ratio of semi-leptonic decays with Cabbibo-allowed and Cabibbo surpressed couplings are given by

$$\frac{g_{cd}^2}{g_{cs}^2} = tan^2\theta_C = \frac{1}{20}$$

The third generation

→ The existence of the c-quark was first predicted from the lepton-quark symmetry.

→ After the discovery of τ, ν_τ, and b, the sixth quark was predicted to complete the symmetry and the top-quark was discovered in 1994 with a mass of about 180 GeV/c^2.

❖ The two generation quark mixing is conveniently written in matrix form as:

$$\begin{pmatrix} d' \\ s' \end{pmatrix} = \begin{pmatrix} \cos\theta_C & \sin\theta_C \\ -\sin\theta_C & \cos\theta_C \end{pmatrix} \begin{pmatrix} d \\ s \end{pmatrix} \tag{105}$$

→ A third generation: gives rise to the Cabibbo-Kobayashi-Maskawa (CKM) matrix $V_{\alpha\beta}$:

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix} \tag{106}$$

→ The coupling constants are then:

$$g_{\alpha\beta} = g_W V_{\alpha\beta} \qquad (\alpha = u, c, t; \beta = d, s, b) \qquad (107)$$

→ If the mixing between the b quark and (d,s) quarks can be neglected, the CKM matrix is reduced to

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \approx \begin{pmatrix} \cos\theta_C & \sin\theta_C & 0 \\ -\sin\theta_C & \cos\theta_C & 0 \\ 0 & 0 & 1 \end{pmatrix} \qquad (108)$$

and hence b'=b.

→ Since the two-generation mixing model agree well with data V_{ub}, V_{cb}, V_{td} and V_{ts} must be small.

→ The b-quark is heavy (mass=4.5 GeV) and it can decay to the lighter u and c quarks as in the following decay:

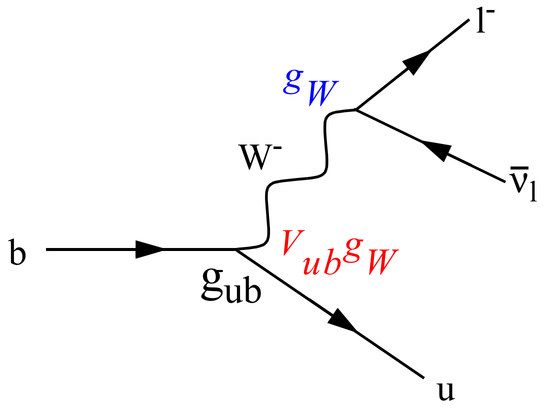


Figure 99: The semi-leptonic decay of the b-quark to a lighter quark.

→ This decay modes has a rate proportional to the squared couplings:

$$|g_{ub}|^2 = |V_{ub}|^2 g_W^2 \qquad (109)$$

→ If V_{ub} and V_{cb} are 0, b-quark must be stable. Experimentally, $\tau_b \approx 10^{-12} s$ which is a long but not infinite lifetime.

❖ The weak b-quark decay can be compared to that of the τ-lepton.

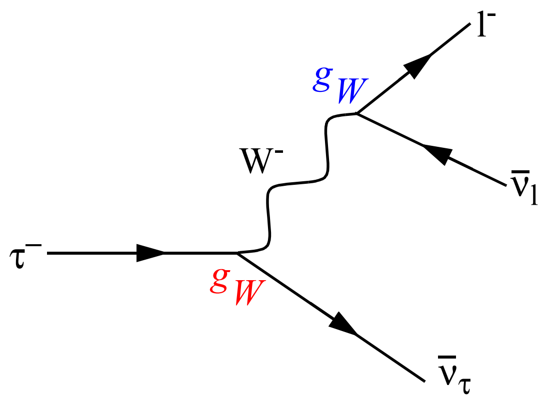


Figure 100: The decay of the τ-lepton.

→ This decay should have decay rates proportional to g_W^2.

If one assumes that $|V_{ub}| = 1$ then one can predict what lifetime this would result in for the b-meson:

$$\tau_b \approx \frac{1}{N}\left(\frac{m_\tau}{m_b}\right)^5 \bullet \tau_\tau \approx 10^{-15} s$$

where $\tau_\tau \approx 3 x 10^{-13}$ s is the lifetime of the τ-lepton, N is the number of possible b-quark decays per analogous τ-decays (3 or 4) and m_τ and m_b are the masses of the τ and the b.

Conclusion:

$$|V_{ub}| = |V_{cb}| = 0 \qquad \tau_b \approx \infty$$

$$|V_{ub}| = 1 \qquad \tau_b \approx 10^{-15}s$$

$$\text{Experimentally:} \qquad \tau_b \approx 10^{-12}s$$

The most precise measurements at present yield

$$|V_{ub}| \approx 0,004 \quad and \quad |V_{cb}| \approx 0,04 \qquad (110)$$

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \approx \begin{pmatrix} \cos\theta_C & \sin\theta_C & 0,004 \\ -\sin\theta_C & \cos\theta_C & 0,04 \\ \sim 0 & \sim 0 & \sim 1 \end{pmatrix}$$

$$\approx \begin{pmatrix} 0,98 & 0,22 & 0,004 \\ -0,22 & 0,98 & 0,04 \\ \sim 0 & \sim 0 & \sim 1 \end{pmatrix}$$

The top-quark

❖ The top-quark is much heavier than even W bosons and can decay by

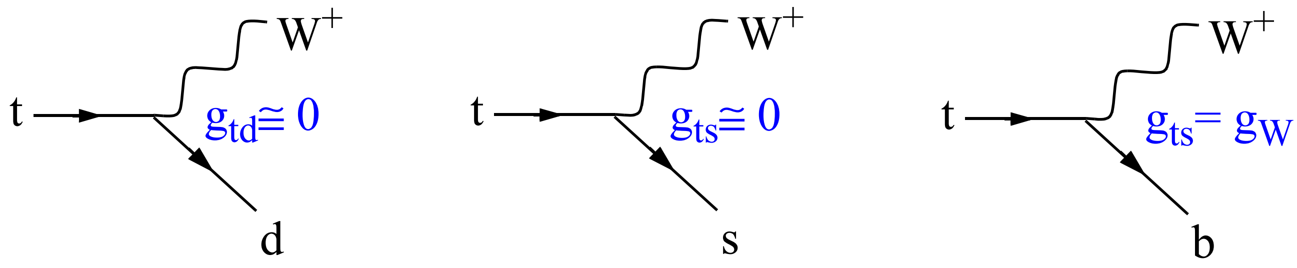


Figure 101: Weak interactions involving the top-quark.

❖ Since $g_{td} \cong 0$ and $g_{ts} \cong 0$ the only significant decay mode of the t-quark is

$$t \to W^{+} + b \qquad (111)$$

with a rate proportional to

$$\alpha_W = g_W^2 / 4\pi \approx 4{,}2\times10^{-3}$$

Estimation of the decay width: $\Gamma \sim \alpha_W m_t \sim 1\ GeV$ suggests a very short lifetime for the top-quark:

$$\tau_t \approx 4\times10^{-25} s \qquad (112)$$

Discovery of the top quark.

❖ The top-quark was discovered in $p\bar{p}$-collisions in the Collider Detector at Fermilab (CDF) in 1994. The Tevatron accelerator with a collision energy = 1.8 TeV was used.

→ PRODUCTION: In proton colliders, pairs of top quarks are produced by the quark-antiquark annihilation process:

$$q + \bar{q} \rightarrow g \rightarrow t + \bar{t} \qquad (113)$$

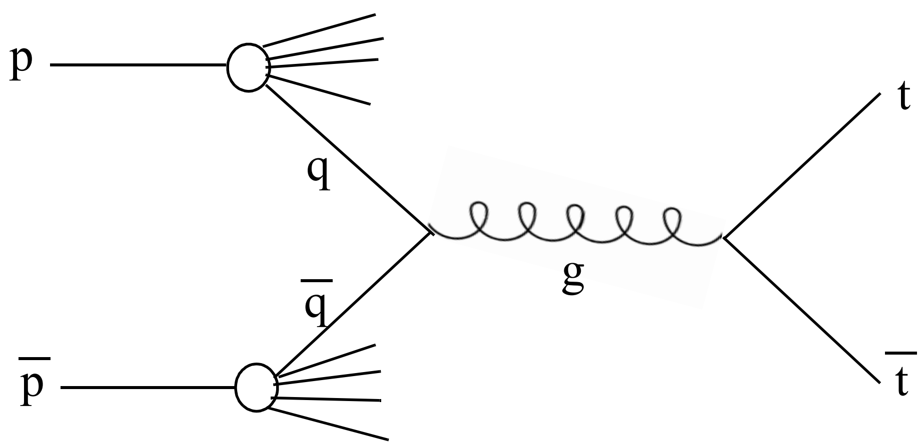


Figure 102: The production of top quarks in $p\bar{p}$ annihilation.

→ DECAY: The top quark decay in most cases to a b-quark and to a W which in turn decays to leptons or quarks:

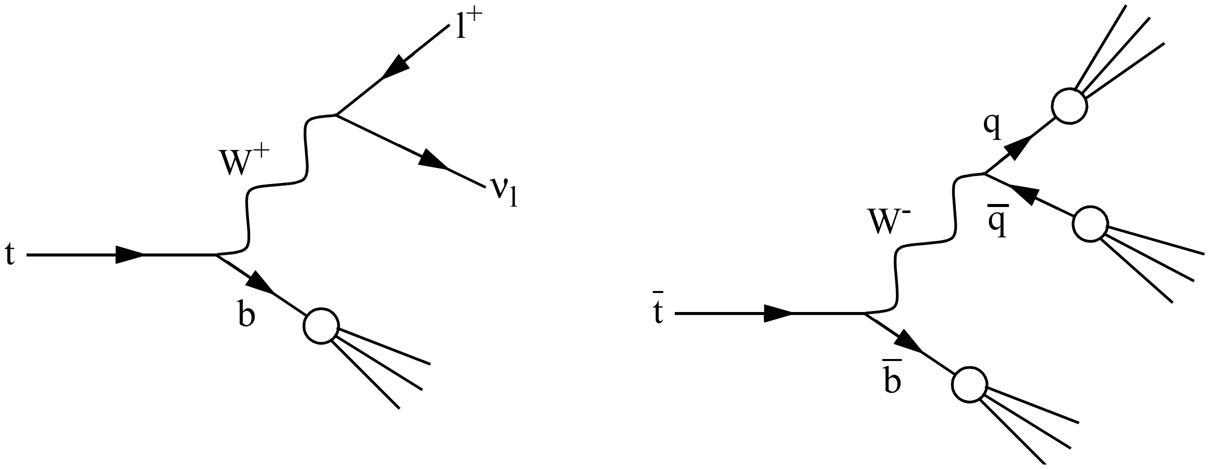


Figure 103: The decay of top quarks.

The final state is a complex mix of jets and leptons:

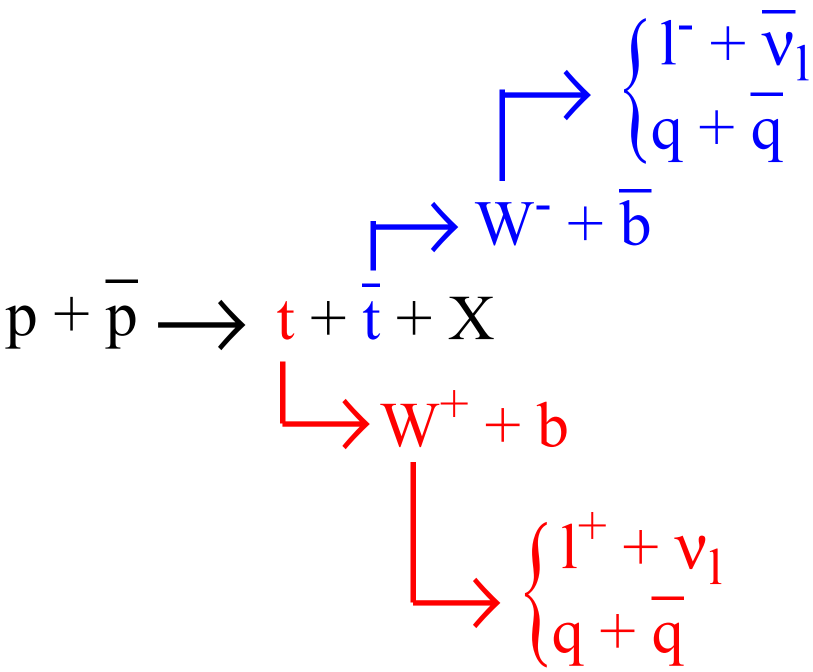
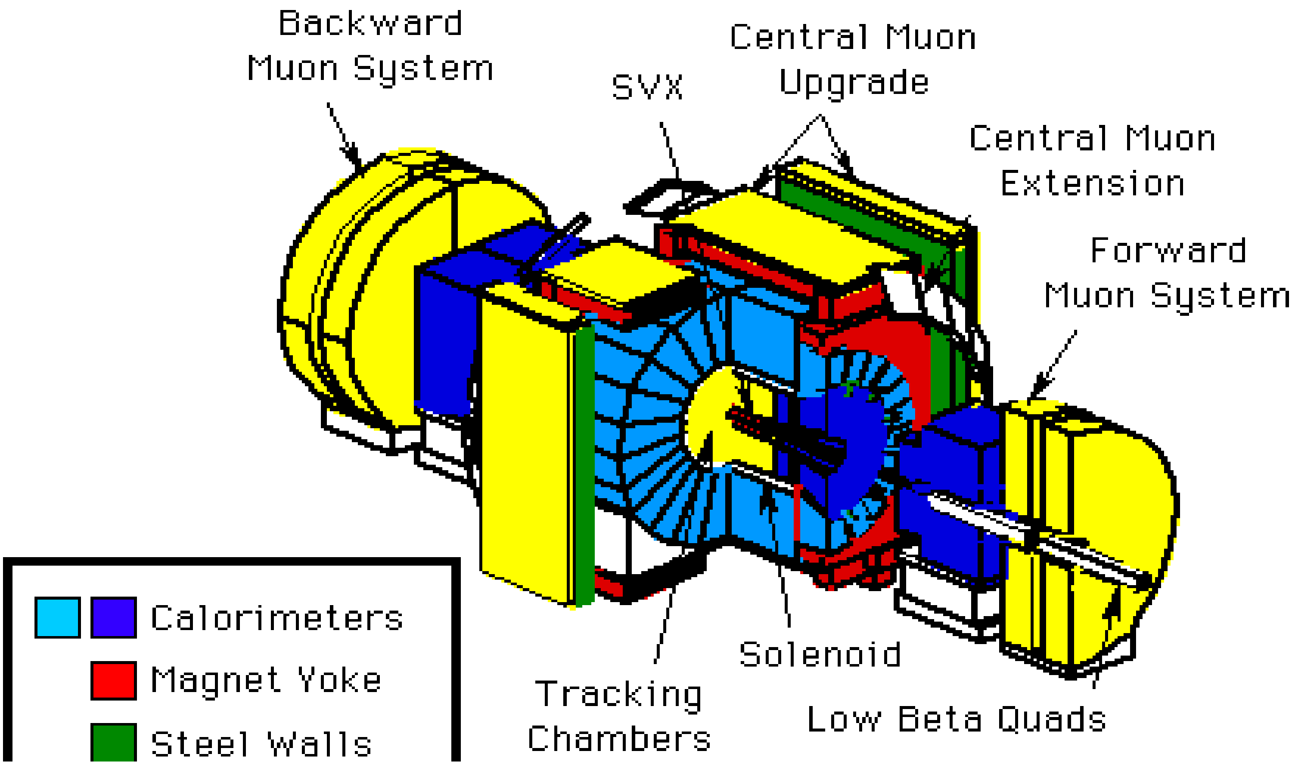


Figure 104: The production of $t\bar{t}$-pairs in $p\bar{p}$-collisions.



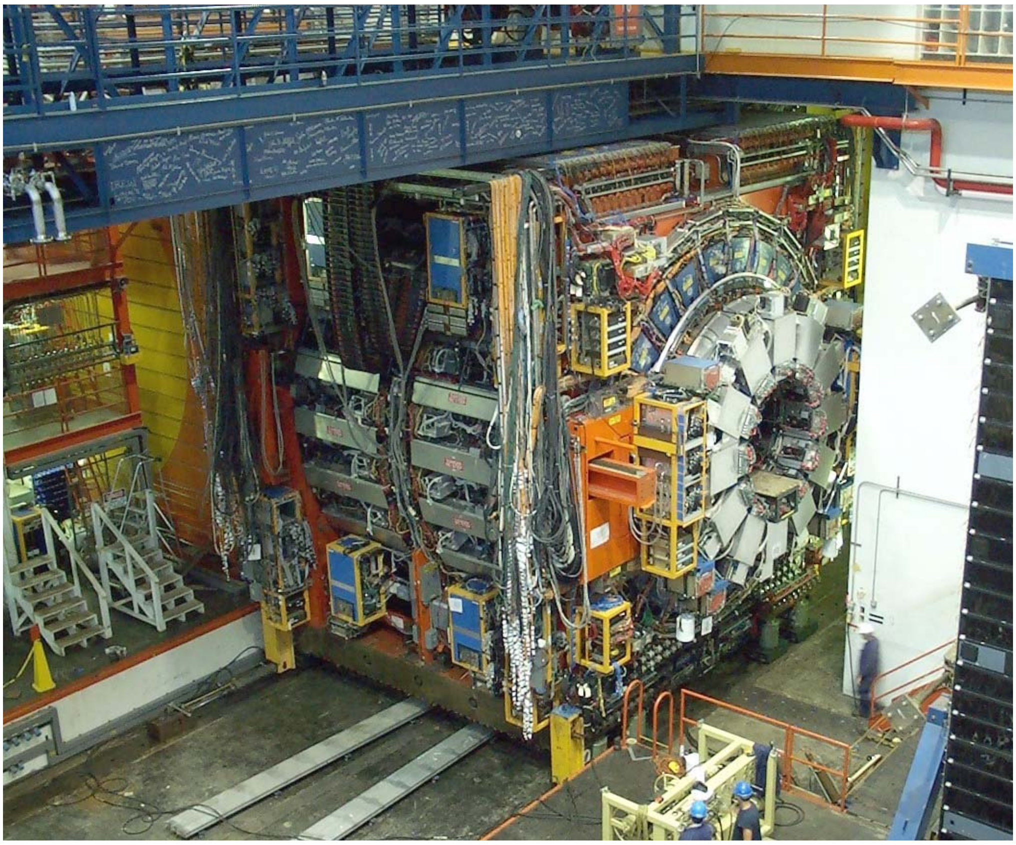


Figure 105: The CDF experiment which discovered the top quark.

→ The events that the CDF collaboration searched for were:

$$t + \bar{t} \rightarrow b + l + \nu + \bar{b} + q + \bar{q}$$

where the lepton l was either an electron or a muon.

→ SIGNATURE:

Two b-jets
Two light quark jets
Missing energy from the neutrino
Lepton with large transverse energy

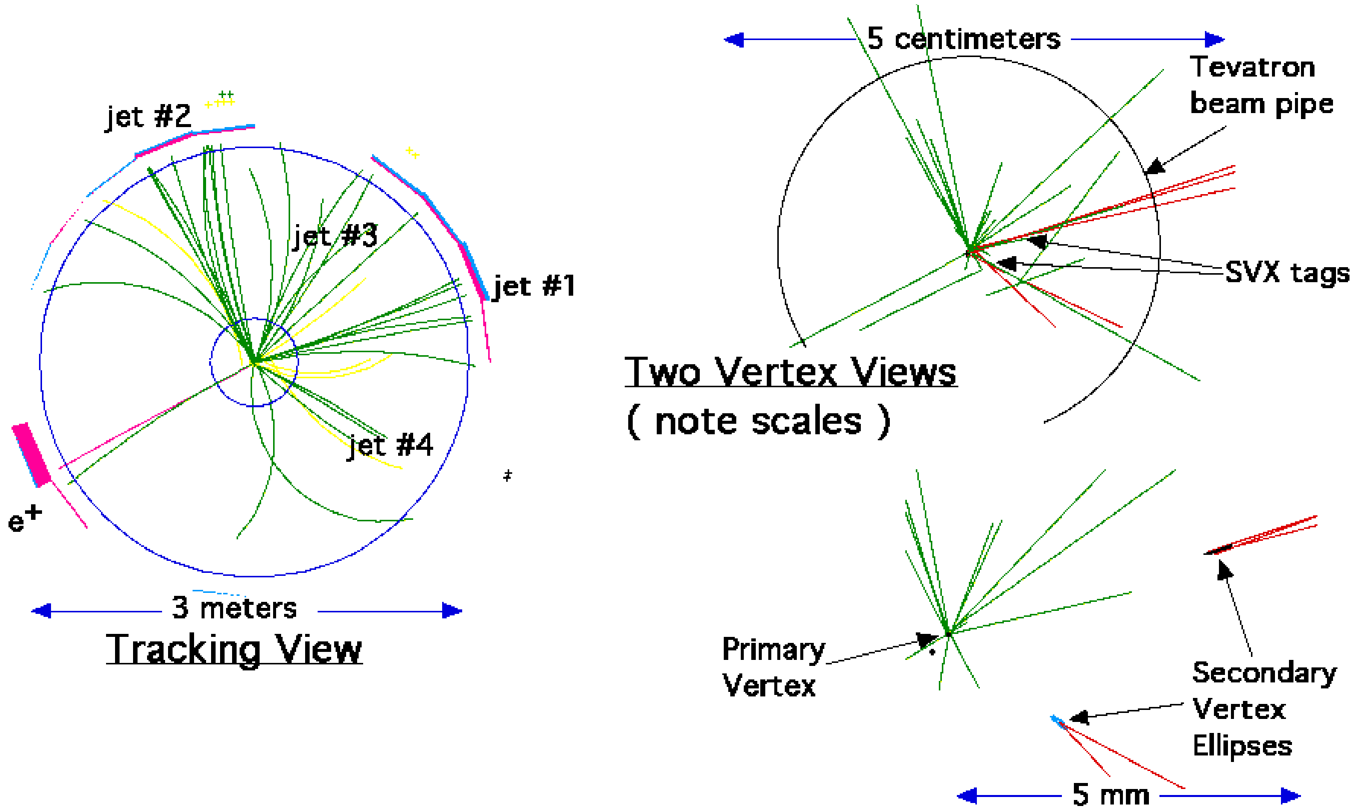


Figure 106: A top event in the CDF experiment.

→ The measured mass distribution (by the CDF experiment) of the events with the required signature was:

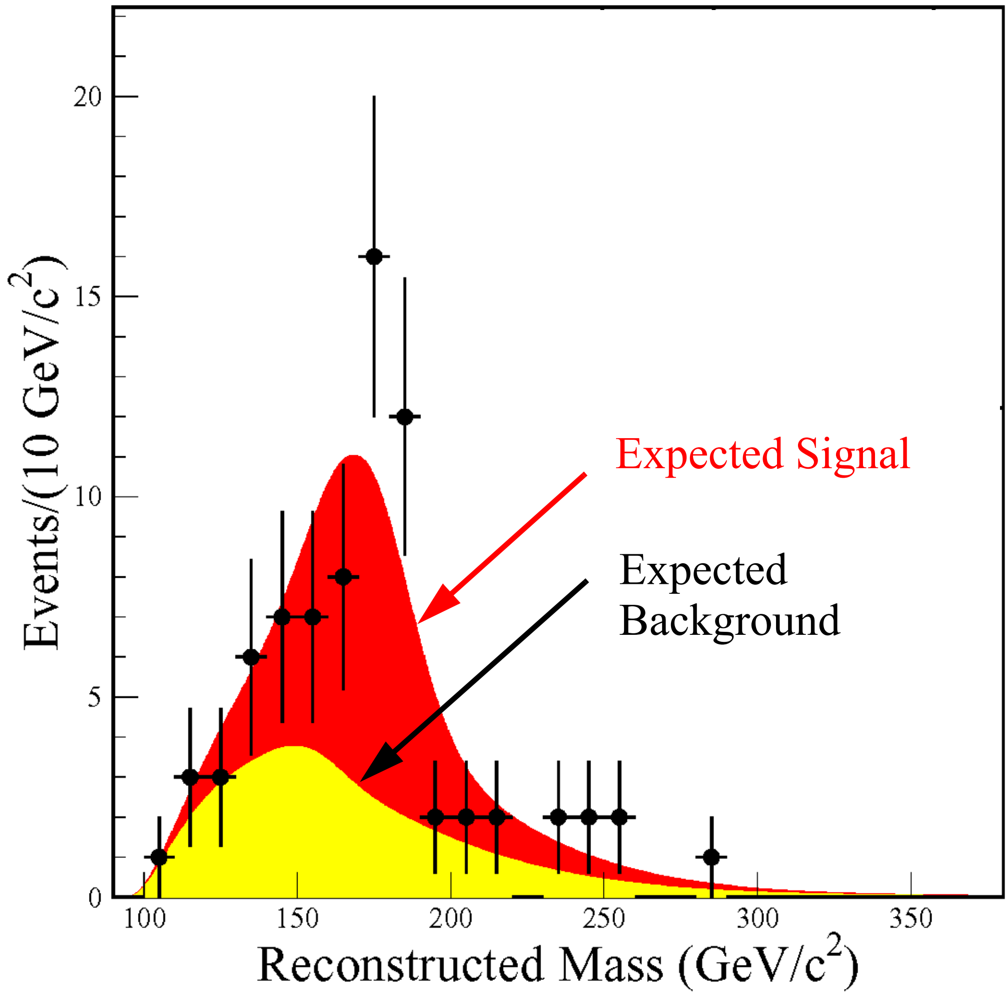


Figure 107: Mass distribution of top events.

→ The top-quark's mass was measured to be

$$M_t = 176 \pm 5\ GeV$$

Summary

• W and Z bosons.

a) W and Z are the force carriers in weak interaction.

b) They are very massive spin-1 bosons.

c) Processes with W exchange are called charged current reactions and processes with Z exchange neutral current reactions.

• Discovery of the W and Z bosons.

d) The W and Z were discovered in $p\bar{p}$-collisions by the UA1 and UA2 experiments at CERN.

e) The W and Z decays to pairs of leptons or quarks.

f) One has measured the W and Z mass with a large accuracy at the LEP accelerator: M_W=80GeV and M_Z=91 GeV.

• Charged current reactions.

g) Leptonic weak interactions can be built from a set of basic vertices.

h) The weak strength parameter is of the same size as that of the electromagnetic interaction.

• Weak interactions of hadrons.

i) There is a lepton-quark symmetry between particles in the same generation.

j) The quark mixing scheme describe how quarks with different flavours interact with each other.

• The third generation.

k) Quark mixing is described by the CKM-matrix.

l) The mixing is much smaller for the heavy quarks than the light quarks.

• The top quark.

m) The top decay in most cases to a W and a b.

n) It was discovered in the CDF experiment.